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# POSITION VECTORS OF SMARANDACHE $\mathcal{D}$ -Tg CURVES OF TIMELIKE BIHARMONIC $\mathcal{D}$ -HELICES ACCORDING TO DARBOUX FRAME ON NON-DEGENERATE TIMELIKE SURFACES IN THE LORENTZIAN HEISENBERG GROUP $\mathbb{H}$

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ABSTRACT. In this paper, we study Smarandache  $\mathcal{D}-\mathbf{Tg}$  curves of timelike biharmonic  $\mathcal{D}$ -helices according to Darboux frame on non-degenerate timelike surfaces in the Lorentzian Heisenberg group  $\mathbb{H}$ . We obtain parametric equation Smarandache  $\mathcal{D}-\mathbf{Tg}$  curves of timelike biharmonic  $\mathcal{D}$ -helices in the Lorentzian Heisenberg group  $\mathbb{H}$ . Moreover, we illustrate the figure of our main theorem.

#### 1. INTRODUCTION

Harmonic maps  $f: (M,g) \longrightarrow (N,h)$  between Riemannian manifolds are the critical points of the energy

(1.1) 
$$E(f) = \frac{1}{2} \int_{M} |df|^2 v_g,$$

and they are therefore the solutions of the corresponding Euler–Lagrange equation. This equation is given by the vanishing of the tension field

(1.2) 
$$\tau(f) = \operatorname{trace} \nabla df.$$

As suggested by Eells and Sampson in [3], we can define the bienergy of a map f by

(1.3) 
$$E_{2}(f) = \frac{1}{2} \int_{M} |\tau(f)|^{2} v_{g},$$

and say that is biharmonic if it is a critical point of the bienergy.

Jiang derived the first and the second variation formula for the bienergy in [7], showing that the Euler-Lagrange equation associated to  $E_2$  is

(1.4) 
$$\tau_2(f) = -\mathcal{J}^f(\tau(f)) = -\Delta \tau(f) - \operatorname{trace} R^N(df, \tau(f)) df$$
$$= 0,$$

where  $\mathcal{J}^{f}$  is the Jacobi operator of f. The equation  $\tau_{2}(f) = 0$  is called the biharmonic equation. Since  $\mathcal{J}^{f}$  is linear, any harmonic map is biharmonic. Therefore, we are interested in proper biharmonic maps, that is non-harmonic biharmonic maps.

In this paper, we study Smarandache  $\mathcal{D}$ -**Tg** curves of timelike biharmonic  $\mathcal{D}$ -helices according to Darboux frame on non-degenerate timelike surfaces in the

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Lorentzian Heisenberg group  $\mathbb{H}$ . We obtain parametric equation Smarandache  $\mathcal{D}$ -**Tg** curves of timelike biharmonic  $\mathcal{D}$ -helices in the Lorentzian Heisenberg group  $\mathbb{H}$ . Moreover, we illustrate the figure of our main theorem.

### 2. Lorentzian Heisenberg Group $\mathbb{H}$

Heisenberg group plays an important role in many branches of mathematics such as representation theory, harmonic analysis, PDEs or even quantum mechanics, where it was initially defined as a group of  $3 \times 3$  matrices

$$\left\{ \left( \begin{array}{ccc} 1 & x & z \\ 0 & 1 & y \\ 0 & 0 & 1 \end{array} \right) : \ x, y, z \in \mathbb{R} \right\}$$

with the usual multiplication rule.

We will use the following complex definition of the Heisenberg group  $\mathbb{H}$ .

$$\mathbb{H} = \mathbb{C} \times \mathbb{R} = \{(w, z) : w \in \mathbb{C}, z \in \mathbb{R}\}$$

with

$$(w,z)*(\overline{w},\overline{z})=(w+\overline{w},z+\overline{z}+\operatorname{Im}(\langle w,\overline{w}
angle)),$$

where  $\langle , \rangle$  is the usual Hermitian product in  $\mathbb{C}$ .

The identity of the group is (0,0,0) and the inverse of (x,y,z) is given by (-x,-y,-z).

Let  $a = (w_1, z_1)$ ,  $b = (w_2, z_2)$  and  $c = (w_3, z_3)$ . The commutator of the elements  $a, b \in \mathbb{H}$  is equal to

$$\begin{aligned} [a,b] &= a * b * a^{-1} * b^{-1} \\ &= (w_1, z_1) * (w_2, z_2) * (-w_1, -z_1) * (-w_2, -z_2) \\ &= (w_1 + w_2 - w_1 - w_2, z_1 + z_2 - z_1 - z_2) \\ &= (0, \alpha) \,, \end{aligned}$$

where  $\alpha \neq 0$  in general. For example

$$[(1,0),(i,0)] = (0,2) \neq (0,0)$$

Which shows that  $\mathbb H$  is not abelian.

On the other hand, for any  $a, b, c \in \mathbb{H}$ , their double commutator is

$$[[a, b], c] = [(0, \alpha), (w_3, z_3)]$$
  
= (0, 0).

This implies that  $\mathbb{H}$  is a nilpotent Lie group with nilpotency 2.

The left-invariant Lorentz metric on  $\mathbb H$  is

$$\rho = -dx^2 + dy^2 + (xdy + dz)^2.$$

The following set of left-invariant vector fields forms an orthonormal basis for the corresponding Lie algebra:

(2.1) 
$$\left\{ \mathbf{e}_1 = \frac{\partial}{\partial z}, \ \mathbf{e}_2 = \frac{\partial}{\partial y} - x \frac{\partial}{\partial z}, \ \mathbf{e}_3 = \frac{\partial}{\partial x} \right\}.$$

The characterising properties of this algebra are the following commutation relations:

(2.2) 
$$\rho(\mathbf{e}_1, \mathbf{e}_1) = \rho(\mathbf{e}_2, \mathbf{e}_2) = 1, \ \rho(\mathbf{e}_3, \mathbf{e}_3) = -1.$$

# 3. Timelike Biharmonic $\mathcal{D}$ -Helices According to Darboux Frame on a Non-Degenerate Timelike Surface in the Lorentzian Heisenberg Group $\mathbb{H}$

Let  $\Pi \subset \mathbb{H}$  be a timelike surface with the unit normal vector  $\mathbf{n}$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . If  $\gamma$  is a timelike curve on  $\Pi \subset \mathbb{H}$ , then we have the Frenet frame  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  and Darboux frame  $\{\mathbf{T}, \mathbf{n}, \mathbf{g}\}$  with spacelike vector  $\mathbf{g} = \mathbf{T} \wedge \mathbf{n}$ along the curve  $\gamma$ . Let  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  be the Frenet frame fields tangent to the Lorentzian Heisenberg group  $\mathbb{H}$  along  $\gamma$  defined as follows:

**T** is the unit vector field  $\gamma'$  tangent to  $\gamma$ , **N** is the unit vector field in the direction of  $\nabla_{\mathbf{T}}\mathbf{T}$  (normal to  $\gamma$ ) and **B** is chosen so that {**T**, **N**, **B**} is a positively oriented orthonormal basis. Then, we have the following Frenet formulas:

(3.1) 
$$\nabla_{\mathbf{T}} \mathbf{T} = \kappa \mathbf{N},$$
$$\nabla_{\mathbf{T}} \mathbf{N} = \kappa \mathbf{T} + \tau \mathbf{B},$$
$$\nabla_{\mathbf{T}} \mathbf{B} = -\tau \mathbf{N},$$

where  $\kappa$  is the curvature of  $\gamma$  and  $\tau$  is its torsion and

(3.2) 
$$\rho(\mathbf{T}, \mathbf{T}) = -1, \ \rho(\mathbf{N}, \mathbf{N}) = 1, \ \rho(\mathbf{B}, \mathbf{B}) = 1,$$
$$\rho(\mathbf{T}, \mathbf{N}) = \rho(\mathbf{T}, \mathbf{B}) = \rho(\mathbf{N}, \mathbf{B}) = 0.$$

Let  $\theta$  be the angle between **N** and **n**. The relationships between  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  and  $\{\mathbf{T}, \mathbf{n}, \mathbf{g}\}$  are as follows:

(3.3) 
$$\mathbf{T} = \mathbf{T},$$
$$\mathbf{N} = \cos \theta \mathbf{n} + \sin \theta \mathbf{g},$$
$$\mathbf{B} = \sin \theta \mathbf{n} - \cos \theta \mathbf{g},$$

and

(3.4) 
$$\begin{aligned} \mathbf{T} = \mathbf{T}, \\ \mathbf{g} = \sin \theta \mathbf{N} - \cos \theta \mathbf{B}, \\ \mathbf{n} = \cos \theta \mathbf{N} + \sin \theta \mathbf{B}. \end{aligned}$$

By differentiating (3.4), using (3.1), (3.3) and Frenet formulas we obtain

(3.5) 
$$\nabla_{\mathbf{T}} \mathbf{T} = (\kappa \cos \theta) \mathbf{n} + (\kappa \sin \theta) \mathbf{g},$$
$$\nabla_{\mathbf{T}} \mathbf{g} = (-\kappa \sin \theta) \mathbf{T} + \left(\tau + \frac{d\theta}{ds}\right) \mathbf{n},$$
$$\nabla_{\mathbf{T}} \mathbf{n} = (-\kappa \cos \theta) \mathbf{T} - \left(\tau + \frac{d\theta}{ds}\right) \mathbf{g}.$$

If we represent  $\kappa \cos \theta$ ,  $\kappa \sin \theta$  and  $\tau + \frac{d\theta}{ds}$  with the symbols  $\kappa_n$ ,  $\kappa_g$ , and  $\tau_g$  respectively, then the equations in (3.5) can be written as

(3.6) 
$$\nabla_{\mathbf{T}} \mathbf{T} = \kappa_g \mathbf{g} + \kappa_n \mathbf{n},$$
$$\nabla_{\mathbf{T}} \mathbf{g} = -\kappa_g \mathbf{T} + \tau_g \mathbf{n},$$
$$\nabla_{\mathbf{T}} \mathbf{n} = -\kappa_n \mathbf{T} - \tau_g \mathbf{g}$$

With respect to the orthonormal basis  $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ , we can write

(3.7) 
$$\mathbf{T} = T_1 \mathbf{e}_1 + T_2 \mathbf{e}_2 + T_3 \mathbf{e}_3, \\ \mathbf{g} = g_1 \mathbf{e}_1 + g_2 \mathbf{e}_2 + g_3 \mathbf{e}_3, \\ \mathbf{n} = n_1 \mathbf{e}_1 + n_2 \mathbf{e}_2 + n_3 \mathbf{e}_3,$$

To separate a curve according to Darboux frame from that of Frenet- Serret frame, in the rest of the paper, we shall use notation for the curve as  $\mathcal{D}$ -curve.

First of all we recall the following well known result (cf. [8]).

**Theorem 3.1.** Let  $\gamma: I \longrightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike curve on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ .  $\gamma$  is a unit speed timelike biharmonic curve on  $\Pi$  if and only if

(3.8) 
$$\kappa_n^2 + \kappa_g^2 = constant \neq 0,$$
$$\kappa_n'' - \kappa_n^3 + \kappa_g \tau_g - \kappa_g^2 \kappa_n + \kappa_g' \tau_g + \kappa_g \tau_g' - \tau_g^2 \kappa_n = \kappa_n (1 - 4g_1^2) - 4\kappa_g n_1 g_1,$$
$$\kappa_g'' - \kappa_g^3 - 2\kappa_n' \tau_g - \kappa_n^2 \kappa_g - \kappa_n \tau_g' - \kappa_g \tau_g^2 = 4\kappa_n n_1 g_1 + \kappa_g (1 - 4n_1^2).$$

**Theorem 3.2.** Let  $\gamma : I \longrightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - helix on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . Then parametric equations of timelike biharmonic  $\mathcal{D}$ - helix are

$$oldsymbol{x}\left(s
ight) = rac{\cosh\mathcal{P}}{\mathcal{M}}(\cosh[\mathcal{M}s]\sinh[\mathcal{N}]+\cosh[\mathcal{N}]\sinh[\mathcal{M}s])+\mathcal{P}_{1},$$

(3.9) 
$$\boldsymbol{y}(s) = \frac{\cosh \mathcal{P}}{\mathcal{M}} (\cosh[\mathcal{N}] \cosh[\mathcal{M}s] + \sinh[\mathcal{N}] \sinh[\mathcal{M}s]) + \mathcal{P}_2,$$

$$\begin{aligned} \boldsymbol{z}\left(s\right) &= \sinh \mathcal{P}s + \frac{\left(\mathcal{N} + \mathcal{M}s\right)}{2\mathcal{M}^{2}} \cosh^{2} \mathcal{P} - \frac{\mathcal{P}_{1}}{\mathcal{M}} \cosh^{2} \mathcal{P} \cosh[\mathcal{N}] \cosh[\mathcal{M}s] \\ &- \frac{\mathcal{P}_{1}}{\mathcal{M}} \cosh^{2} \mathcal{P} \sinh[\mathcal{N}] \sinh[\mathcal{M}s] - \frac{1}{4\mathcal{M}^{2}} \cosh^{2} \mathcal{P} \sinh 2[\mathcal{M}s + \mathcal{N}] + \mathcal{P}_{3}, \end{aligned}$$

where  $\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3$  are constants of integration and

$$\mathcal{M} = \operatorname{sech} \mathcal{P}_{\sqrt{\kappa_n^2 + \kappa_g^2}} - 2 \operatorname{sinh} \mathcal{P}.$$

Using Mathematica in above system we have Fig.1:



**Corollary 3.3.** Let  $\gamma : I \longrightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - curve on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ . Then,

 $\kappa = constant \neq 0.$ 

# 4. Smarandache $\mathcal{D}$ -Tg curves According to Darboux Frame on a Non-Degenerate Timelike Surface in the Lorentzian Heisenberg Group $\mathbb{H}$

**Definition 4.1.** Let  $\gamma : I \longrightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - curve on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ and  $\{\mathbf{T}, \mathbf{n}, \mathbf{g}\}$  be its moving Darboux frame. Smarandache  $\mathcal{D}$ -**Tg** curves are defined by

(4.1) 
$$\tilde{\gamma} = \frac{1}{\kappa_n + \tau_g} \left( \mathbf{T} + \mathbf{g} \right).$$

**Theorem 4.2.** Let  $\gamma : I \longrightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - curve on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ and  $\tilde{\gamma}$  its Smarandache  $\mathcal{D}$ -**T**g curve. Then position vector of Smarandache  $\mathcal{D}$ -**T**g curve is given by

$$\tilde{\gamma} = \left[\frac{\sinh\mathcal{P}}{\kappa_n + \tau_g} - \frac{\cos^2\theta}{\kappa_n (\kappa_n + \tau_g)} \left[\cosh\mathcal{P}\sinh\mathcal{M}s + \mathcal{N}\right]\cosh\mathcal{P}\sinh\left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh\mathcal{P}\right)\right]$$

$$(4.2)$$

$$-\cosh\mathcal{P}\cosh\left[\mathcal{M}s + \mathcal{N}\right]\cosh\mathcal{P}\cosh\left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh\mathcal{P}\right)\right]$$

$$= 1$$

$$+ \left[\frac{\cosh\mathcal{P}\sinh\left[\mathcal{M}s + \mathcal{N}\right]}{\kappa_n + \tau_g} + \frac{1}{\kappa_n + \tau_g} \left[\frac{\cos\theta\sin\theta}{\kappa_n}\cosh\mathcal{P}\cosh\left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh\mathcal{P}\right)\right]$$

$$+ \frac{\cos^2\theta}{\kappa_n}\sinh\mathcal{P}\cosh\mathcal{P}\sinh\left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh\mathcal{P}\right)\right]$$

$$= 1$$

$$+ \left[\frac{\cosh\mathcal{P}\cosh\left[\mathcal{M}s + \mathcal{N}\right]}{\kappa_n + \tau_g} + \frac{1}{\kappa_n + \tau_g} \left[\frac{\cos\theta\sin\theta}{\kappa_n}\cosh\mathcal{P}\sinh\left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh\mathcal{P}\right)$$

$$- \frac{\cos^2\theta}{\kappa_n}\sinh\mathcal{P}\cosh\mathcal{P}\cosh\left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh\mathcal{P}\right)\right]$$

$$= 0$$

where  $\mathcal{P}_1, \mathcal{P}_2, \mathcal{P}_3$  are constants of integration and

$$\mathcal{M} = \operatorname{sech} \mathcal{P}_{\sqrt{\kappa_n^2 + \kappa_g^2 - 2 \sinh \mathcal{P}}}.$$

**Proof.** From the assumption we get

(4.3) 
$$\mathbf{T} = \sinh \mathcal{P} \mathbf{e}_1 + \cosh \mathcal{P} \sinh \left[\mathcal{M} s + \mathcal{N}\right] \mathbf{e}_2 + \cosh \mathcal{P} \cosh \left[\mathcal{M} s + \mathcal{N}\right] \mathbf{e}_3.$$
  
The covariant derivative of the vector field **T** is:

$$\nabla_{\mathbf{T}} \mathbf{T} = \cosh \mathcal{P} \cosh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right) \mathbf{e}_{2} + \cosh \mathcal{P} \sinh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right) \mathbf{e}_{3}.$$

From the above equation, it is seen that

$$\mathbf{g} = -\frac{\cos^2 \theta}{\kappa_n} [\cosh \mathcal{P} \sinh \mathcal{M}s + \mathcal{N}] \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] (\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P}) - \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] (\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P})] \mathbf{e}_1 (4.4) + [\frac{\cos \theta \sin \theta}{\kappa_n} \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] (\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P}) + \frac{\cos^2 \theta}{\kappa_n} \sinh \mathcal{P} \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] (\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P})] \mathbf{e}_2 [\frac{\cos \theta \sin \theta}{\kappa_n} \cosh \mathcal{P} \sinh [\mathcal{M}s + \mathcal{N}] (\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P})] \mathbf{e}_2 - \frac{\cos^2 \theta}{\kappa_n} \sinh \mathcal{P} \cosh \mathcal{P} \cosh [\mathcal{M}s + \mathcal{N}] (\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P})] \mathbf{e}_3.$$

Substituting Eq.(4.4) and Eq.(4.3) into Eq.(4.1) , we obtain Eq.(4.2). This completes the proof.

**Theorem 4.3.** Let  $\gamma : I \longrightarrow \Pi \subset \mathbb{H}$  be a non-geodesic unit speed timelike biharmonic  $\mathcal{D}$ - curve on timelike surface  $\Pi$  in the Lorentzian Heisenberg group  $\mathbb{H}$ 

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and  $\tilde{\gamma}$  its Smarandache  $\mathcal{D}$ -**Tg** curve. Then parametric equations of Smarandache  $\mathcal{D}$ -**Tg** curve are given by

$$x_{\tilde{\gamma}}(s) = \left[\frac{\cosh \mathcal{P} \cosh \left[\mathcal{M} s + \mathcal{N}\right]}{\kappa_n + \tau_g} + \frac{1}{\kappa_n + \tau_g} \left[\frac{\cos \theta \sin \theta}{\kappa_n} \cosh \mathcal{P} \sinh \left[\mathcal{M} s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P}\right) - \frac{\cos^2 \theta}{\kappa_n} \sinh \mathcal{P} \cosh \mathcal{P} \cosh \left[\mathcal{M} s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2 \sinh \mathcal{P}\right)\right],$$

$$y_{\tilde{\gamma}}(s) = \left[\frac{\cosh \mathcal{P} \sinh \left[\mathcal{M}s + \mathcal{N}\right]}{\kappa_n + \tau_g} + \frac{1}{\kappa_n + \tau_g} \left[\frac{\cos \theta \sin \theta}{\kappa_n} \cosh \mathcal{P} \cosh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right) + \frac{\cos^2 \theta}{\kappa_n} \sinh \mathcal{P} \cosh \mathcal{P} \sinh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right)\right],$$
  
$$z_{\tilde{\gamma}}(s) = \left[\frac{\sinh \mathcal{P}}{\kappa_n} - \frac{\cos^2 \theta}{\kappa_n} \left[\cosh \mathcal{P} \sinh \mathcal{M}s + \mathcal{N}\right] \cosh \mathcal{P} \sinh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right)\right],$$

$$z_{\tilde{\gamma}}(s) = \left[\frac{\sin \eta}{\kappa_n + \tau_g} - \frac{\cos \theta}{\kappa_n (\kappa_n + \tau_g)} \left[\cosh \mathcal{P} \sinh \mathcal{M}s + \mathcal{N}\right] \cosh \mathcal{P} \sinh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right) - \cosh \mathcal{P} \cosh \left[\mathcal{M}s + \mathcal{N}\right] \cosh \mathcal{P} \cosh \left[\mathcal{M}s + \mathcal{N}\right] \left(\frac{1}{\mathcal{M}} + 2\sinh \mathcal{P}\right)\right]$$

$$-\left[\frac{\cosh\mathcal{P}\sinh\left[\mathcal{M}s+\mathcal{N}\right]}{\kappa_{n}+\tau_{g}}+\frac{1}{\kappa_{n}+\tau_{g}}\left[\frac{\cos\theta\sin\theta}{\kappa_{n}}\cosh\mathcal{P}\cosh\left[\mathcal{M}s+\mathcal{N}\right]\left(\frac{1}{\mathcal{M}}+2\sinh\mathcal{P}\right)\right.\\\left.+\frac{\cos^{2}\theta}{\kappa_{n}}\sinh\mathcal{P}\cosh\mathcal{P}\sinh\left[\mathcal{M}s+\mathcal{N}\right]\left(\frac{1}{\mathcal{M}}+2\sinh\mathcal{P}\right)\right]\right]\\\left[\frac{\cosh\mathcal{P}\cosh\left[\mathcal{M}s+\mathcal{N}\right]}{\kappa_{n}+\tau_{g}}+\frac{1}{\kappa_{n}+\tau_{g}}\left[\frac{\cos\theta\sin\theta}{\kappa_{n}}\cosh\mathcal{P}\sinh\left[\mathcal{M}s+\mathcal{N}\right]\left(\frac{1}{\mathcal{M}}+2\sinh\mathcal{P}\right)\right.\\\left.-\frac{\cos^{2}\theta}{\kappa_{n}}\sinh\mathcal{P}\cosh\mathcal{P}\cosh\left[\mathcal{M}s+\mathcal{N}\right]\left(\frac{1}{\mathcal{M}}+2\sinh\mathcal{P}\right)\right]\right],$$

where  $\mathcal{N}$  is constant of integration and

$$\mathcal{M} = \operatorname{sech} \mathcal{P}_{\sqrt{\kappa_n^2 + \kappa_g^2}} - 2 \operatorname{sinh} \mathcal{P}.$$

Proof. Omitted.

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